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Trends and Analysis of Power Distribution in VLSI Interconnects

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Abstract

The analysis of effects induced by interconnects become increasingly important as the scale of process technologies steadily shrinks. While most analyses focus on the timing aspects of interconnects, power consumption is also important. We study the trends of interconnect power consumption based on current and future technology node parameters. We also study the analysis method based on a reduced-order model. The relation between power consumption and the poles and residues of a transfer function is derived. The theoretical results can be used for any kind of linear circuits including RLC circuits.

1 Introduction

As the scale of process technologies steadily shrinks and the size of designs increases, interconnects have increasing impact on the area, delay, and power consumption of circuits. Reduction in scale causes a continual increase in interconnect delays, although of course overall circuit performance continues to increase. As regards power, the situation is similar in that the portion of power associated with interconnects is increasing. This is an important fact because the conventional design, analysis, and synthesis of VLSI circuits are based on the assumption that gates are the main sources of on-chip power consumption. In this paper, we study the trends of interconnect power distribution based on technology node parameters extrapolated from ITRS [1] roadmap. The study shows that for a case of optimally buffered interconnect systems about 40% of the total power is consumed at transistors while the remaining power is changed to heat in interconnects.

For analysis of the interconnects, extensive studies have been made of the use of model order reduction over the last few years, following the introduction of Asymptotic Waveform Evaluation [2]. Model order reduction is based on approximating the Laplace-domain transfer function of a linear (or linearized) network by a relatively small number of dominant poles and zeros. Such reduced-order models can be used to predict the time-domain or frequency-domain response of the linear network. Although there has been significant progress in the analysis and simulation of performance-related aspects of VLSI interconnects, less work has been devoted to the analysis of power consumption (or distribution) of interconnects. Furthermore, the analysis of power-related aspects of interconnects is limited to power distribution networks, and deals with quantities such as IR drop, ground bounce, and electromigration. In this paper, we introduce a method based on a reduced-order model that allows the power distribution of interconnects to be analyzed. We show that the power, which inherently involves improper integration, can be derived from the poles and residues of the transfer function, which requires only algebraic computation.



Figure 1. Buffered global interconnects.

2 Trends of Interconnect Power Distribution

In order to investigate power consumption of interconnects, we consider a buffered global interconnect system as shown in Figure 1. Especially, we consider *optimally buffered inter-connect*, where buffer size and interconnect length are determined in such a way that delay is minimized. The optimal buffer size and the optimal interconnect length can be derived from the delay equation of buffered interconnect [3]:

$$t = k \left[p_1 \frac{r_i L}{k} \frac{c_i L}{k} + p_2 \left(\frac{r_t}{h} h c_t + \frac{r_t}{h} h c_j + \frac{r_t}{h} \frac{c_i L}{k} + \frac{r_i L}{k} h c_t \right) \right], \tag{1}$$

where k is the number of interconnect sections consisting of k - 1 buffers, L is the total length of the interconnect, c_i and r_i are capacitance and resistance of the wire of unit length respectively, and c_t , c_j , and r_t are gate capacitance, drain junction capacitance, and drain effective resistance of minimum size MOSFET respectively. The constants p_1 and p_2 depend on the switching model of the buffer, and are about 0.377 and 0.693, respectively, when 50% of the swing at the receiver side is of interest. By differentiating (1) in terms of h, setting the derivative equal to zero, and solving the equation for h, we obtain the optimal buffer size:

$$h_o = \sqrt{\frac{c_i r_t}{r_i r_t}}.$$
(2)

Differentiating (1) again in terms of k, setting the derivative equal to zero, and solving the equation for k gives us the optimal number of stages for the interconnect of length L, denoted by k_o . Dividing L by k_o can be shown to give us the optimal interconnect length of each section:

$$l_{o} = \sqrt{\frac{p_{2}}{p_{1}}} \sqrt{\frac{r_{t}(c_{t} + c_{j})}{r_{i}c_{i}}}.$$
(3)

Technology	V _{dd}	Tox	r_t	c_t
(nm)	(V)	(Å)	$(k\Omega)$	(fF)
180	1.5	25	19	0.443
130	1.5	19	26	0.260
100	1.2	15	27	0.194
70	1.2	12	39	0.119
50	1.1	8	50	0.090

Table 1. Technology parameters

Table 2. Parameters of optimally buffered interconnect systems

Technology	h_o	l_o	R_i	C_i
(nm)		(mm)	(Ω)	(fF)
180	281	2.63	95	176
130	296	1.67	100	109
100	254	1.10	114	70
70	245	0.67	113	41
50	215	0.46	150	27

To obtain the optimal buffer size and the optimal interconnect length for the current and future technology generation, we use technology parameters extrapolated from ITRS [1] roadmap, which are summarized in Table 1, together with (2) and (3). For each technology node, we obtain optimum size of buffer and optimum length of interconnect, which are tabulated in Table 2 together with total resistance (R_i) and capacitance (C_i) of the interconnect.

For each technology node, we configure the circuit shown in Figure 1 with parameters in Table 2. The interconnect is approximated by 5 sections of π -ladder circuits [4]. We obtain the power consumption of the buffer, denoted as P_b , and of the interconnect (sum of power



Figure 2. Trends of power distribution of optimally buffered interconnect system when (a) step and (b) clock is applied at the input of the buffer.

consumption of 5 resistors in π -ladders), denoted as P_i , with SPICE simulation. The ratio of power consumed by the buffer to the total power consumption is defined by

$$\eta_b = \frac{P_b}{P_b + P_i},\tag{4}$$

and similarly for η_i . When the step is applied at the input of the buffer, the ratios of power are shown graphically in Figure 2(a). Since we use the step at falling edge, most of power consumed by the buffer is due to PMOS operated in a linear region, thus we neglect the power consumed by NMOS, which consists of mostly leakage power. The trends show that about 40% of the total power is consumed at the transistor and the remaining power is changed to heat in the interconnect.

The trends do not change significantly when we use clock instead of step and taking all power components of the buffer into account as shown in Figure 2(b). Since clock trees consume significant power and large buffer and long interconnect (even larger and longer than those for

optimally buffered interconnect as presented in this paper) are frequently used to build clock trees [5], the power consumption of interconnects should be considered an important factor in the design of clock distribution networks.

3 Analysis of Interconnect Power Distribution

For a given linear or linearized circuit, the total power consumption is readily obtained. However, this does not give how the power consumption is distributed over circuit elements. In order to find the power consumption (or energy dissipation)¹ of a particular resistor element, we first obtain the reduced-order model of current flowing through the resistor, denoted by $\hat{J}(s)$ (with the corresponding time-domain function $\hat{j}(t)$), using a model order reduction techniques [2, 6]. The approximate energy dissipated by R_i , denoted by \hat{E}_i , during time period $[\tau_1, \tau_2]$ is then given by

$$\hat{E}_i = R_i \int_{\tau_1}^{\tau_2} \hat{j}^2(t) \, dt.$$
(5)

If we are interested in the total energy dissipated by a specific resistor element during signal transition, we can choose to consider a semi-infinite interval of τ , without loss of generality. We make τ_1 the time origin and τ_2 infinite time. Then $\hat{j}(t)$ will reach a steady state, provided that $\hat{j}(t)$ corresponds to the reduced-order model of an individual transition. This leads us to the improper integral

$$\hat{E}_i = R_i \int_0^\infty \hat{j}^2(t) dt.$$
(6)

The direct computation of improper integration in (6) is difficult, especially if there are multipleorder poles or if $\hat{j}^2(t)$ is expressed as combination of functions other than exponentials. Fortunately, we can avoid this by deriving a general relation between improper integration in the time-domain and algebraic computation in the *s*-plane as expressed by the following theorem and proved in [7].

¹*Power consumption* and *energy dissipation* are used interchangeably. More precisely, power consumption in this paper means *average power consumption*, which is equal to energy dissipation divided by the time period of interest.

Theorem 1 If the Laplace transform of a time-domain signal h(t), denoted by H(s), has q singularities in the left half of the s-plane, then

$$\int_{0}^{\infty} h^{2}(t) dt = \sum_{i=1}^{q} \tilde{r}_{i},$$
(7)

where \tilde{r}_i is a residue of H(-s)H(s) at the singularity of H(s).

Note that the only constraint imposed by Theorem 1 is that the transfer function has singularities to the left of the *s*-plane, which is a typical situation because we are concerned mostly with stable systems. Thus, we can apply the same relation even for RLC circuits. If we have a reduced-order model of H(s), then \tilde{r}_i can be obtained by a matrix computation involving the moments of $\hat{H}(-s)\hat{H}(s)$ and the singularities of $\hat{H}(s)$ [2]. In the case when all the singularities are simple poles, we obtain the less complicated relation expressed by the following theorem.

Theorem 2 If the Laplace transform of a time-domain signal h(t), denoted by H(s), has q simple poles in the left half of s-plane, then

$$\int_0^\infty h^2(t) dt = \sum_{i=1}^q r_i H(-p_i),$$
(8)

where r_i is a residue of H(s) at the pole p_i of H(s).

As an example, suppose we have

$$H(s) = \frac{3s+5}{(s+1)(s+2)} = \frac{2}{s+1} + \frac{1}{s+2} = \frac{r_1}{s-p_1} + \frac{r_2}{s-p_2}.$$
(9)

We apply the above theorem:

$$r_1H(-p_1) + r_2H(-p_2) = 2\left(\frac{2}{1+1} + \frac{1}{1+2}\right) + 1\left(\frac{2}{2+1} + \frac{1}{2+2}\right) = \frac{43}{12}.$$
 (10)

It can be easily shown that this is equal to the result of direct computation of improper integration:

$$\int_0^\infty h^2(t) dt = \int_0^\infty (2e^{-t} + e^{-2t})^2 dt.$$
(11)



Figure 3. Comparison of the energy distribution for a randomly generated circuit with 300 nodes.

To summarize, the overall procedure to analyze power distribution of linear circuit (either RC or RLC circuit) is: at each resistor element, we obtain the reduced-order model of current flowing through the resistor $(\hat{J}(s))$; we apply either Theorem 1 or 2 to compute the improper integration, which is then multiplied by the resistance to obtain the energy dissipation.

In order to verify the validity of the proposed analysis method, we implement a prototype tool written in C^{++} , and based on the results presented in this section with moment matchingbased model order reduction [2]. The program reads in a circuit in a SPICE-like format and outputs the power distribution of the interconnect. For the experiments, we randomly generate RC tree networks while varying the number of nodes from 100 to 500, and compare the energy distribution obtained by SPICE with that obtained by our method. As an example, Figure 3 shows the result for a circuit with 300 nodes. The result shows that a single pole approximation is quite accurate for most of cases. To understand this, first note that the area under the current waveform when it is approximated by a single pole is equal to that under the exact waveform. Since the exact waveform is bell-shaped (except for the driving end) while the approximated one decays monotonically, the accuracy of energy approximation with a single pole depends on the peakness exhibited by the curve, because we are interested in the area under the square of the waveform. If the peakness occurs highly skewed to the left side, meaning that the waveform has high frequency components, the approximated current waveform has a large error around its peak (although the area underneath is correct) and this error becomes more significant when we compute the square of the current waveform, as we must. However, the number of waveforms with such dominant high frequency components is not large for most of practical circuits. The approximation with two poles gives accurate result, which is consistently observed with other circuits.

4 Conclusion

We study the trends of interconnect power distribution based on current and future technology node parameters. The study shows that for a case of optimally buffered interconnect systems about 60% power is changed to heat in interconnects.

We describe a method for the power distribution analysis of an interconnect based on a reduced-order model. We show that power consumption can be computed efficiently in the *s*-domain using an algebraic formulation, instead of improper integration in the time-domain. The theoretical results relies on the poles and residues of a transfer function, and can thus be used in any kind of model order reduction technique.

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